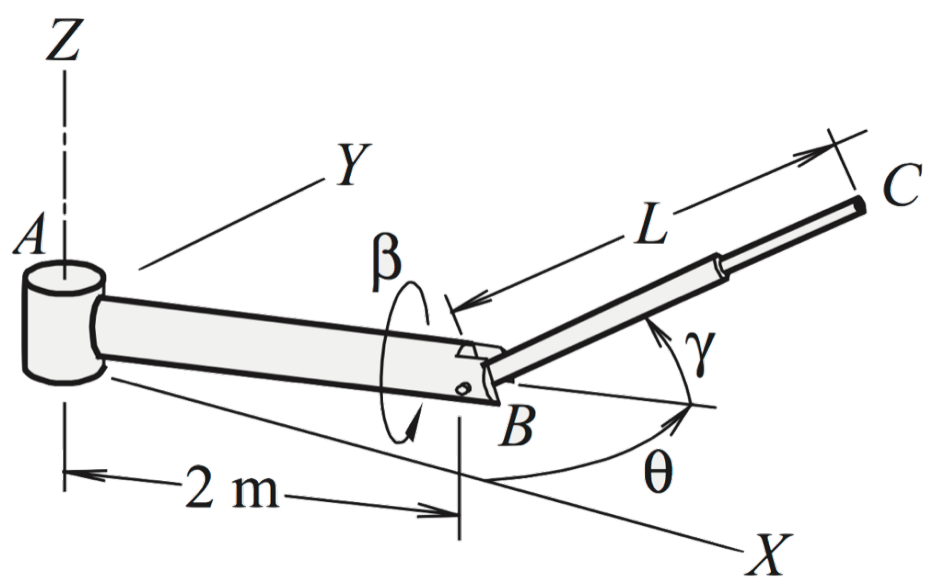


### Example 3.6



AB rotates by  $\theta$  about Z

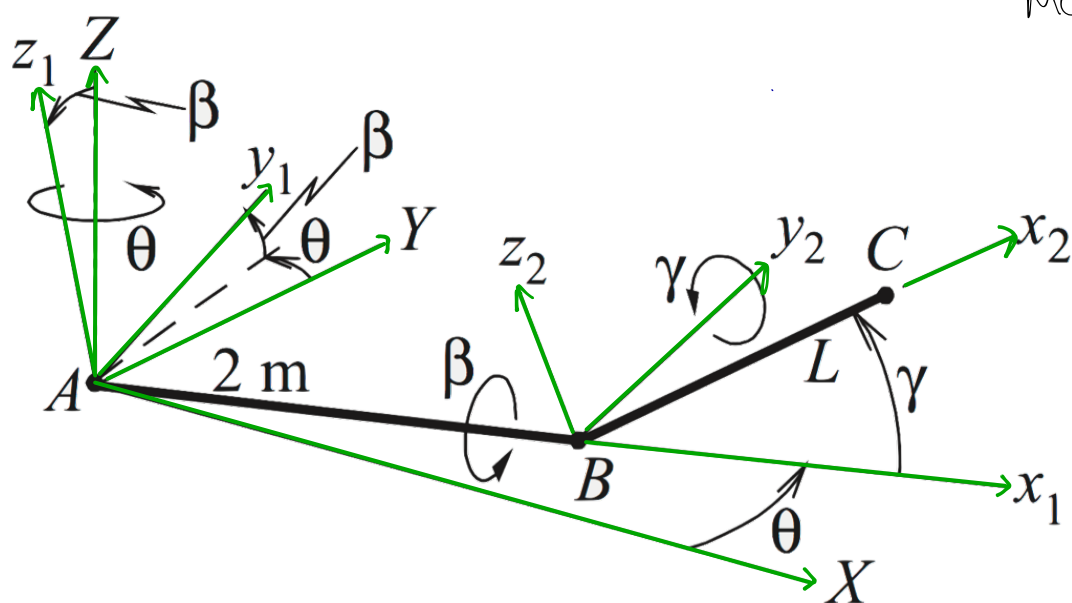
BC rotates by  $\beta$  about AB

BC rotates by  $\gamma$  about pin at B

Given:  $\theta = 50^\circ$ ,  $\beta = 30^\circ$ ,  $\gamma = 60^\circ$

$L_0 = 0.5 \text{ m} \rightarrow L_f = 1.5 \text{ m}$

What is displacement of C  $\Delta \vec{r}_C$  from these moves?



Define  $x_1, y_1, z_1$  fixed to AB

$x_2, y_2, z_2$  fixed to BC

$z_1 = k$  when  $\beta = 0$   
and

$y_2$  is aligned with pin at B

We'll solve by solving for displacement of B in  $x_1, y_1, z_1$ , then for C in  $x_2, y_2, z_2$

$$\Delta \vec{r}_B = \Delta \vec{r}_A + R_{1f} (\Delta \vec{r}_B)_{x_1 y_1 z_1} + (R_{1f} - R_{10}) \begin{bmatrix} x_{1B} \\ y_{1B} \\ z_{1B} \end{bmatrix}_0$$

$$\Delta \vec{r}_C = \Delta \vec{r}_B + R_{2f} (\Delta \vec{r}_C)_{x_2 y_2 z_2} + (R_{2f} - R_{20}) \begin{bmatrix} x_{2C} \\ y_{2C} \\ z_{2C} \end{bmatrix}_0$$

Initially, there is no rotation and  $x_1, y_1, z_1$ ,  $x_2, y_2, z_2$ , and  $X, Y, Z$  are aligned, so

$$R_{10} = R_{20} = I \quad (\text{a n base } = \text{ still Identity matrix})$$

### Example 3.6 (cont.)

Q: What is the rotation matrix from XYZ to  $x_1y_1z_1$

① Rotation of  $\theta$  about Z

② Rotation of  $\beta$  about  $x_1$

$$R_{1f} = R_{\beta} R_{\theta} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\beta & \sin\beta \\ 0 & -\sin\beta & \cos\beta \end{bmatrix} \begin{bmatrix} \cos\theta & \sin\theta & 0 \\ -\sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Q: Rotation of  $x_2y_2z_2$  relative to XYZ?

Just a rotation of  $-\gamma$  about  $y_2$  relative to  $x_1y_1z_1$ , so

$$R_{2f} = R_{\gamma} R_{\beta} R_{\theta} = \begin{bmatrix} \cos\gamma & 0 & \sin\gamma \\ 0 & 1 & 0 \\ -\sin\gamma & 0 & \cos\gamma \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\beta & \sin\beta \\ 0 & -\sin\beta & \cos\beta \end{bmatrix} \begin{bmatrix} \cos\theta & \sin\theta & 0 \\ -\sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Now, all that's left are the relative position + displacements

$$\Delta \vec{r}_A = 0$$

$$(\vec{r}_{B/A})_0 = 2\vec{e}_1 \leftarrow B \text{ is } 2\text{m along } \vec{e}_1 \text{ from } A$$

$$(\Delta \vec{r}_B)_{x_1y_1z_1} = 0 \leftarrow \text{It doesn't change position in the } x_1y_1z_1 \text{ frame}$$

$$(\vec{r}_{C/B})_0 = 0.5\text{m } \vec{e}_2 \leftarrow C \text{ starts } 0.5\text{m from } B \text{ along } \vec{e}_2$$

$$(\Delta \vec{r}_C)_{x_2y_2z_2} = 1\text{m } \vec{e}_2 \leftarrow C \text{ moves from } 0.5\text{m to } 1.5\text{m as rotating}$$

## Time Derivatives (Sec. 3.3)

Now, look at an infinitesimally small displacement  $\rightarrow$  infinitesimal small amount of time

So

$$\Delta \bar{r}_p \rightarrow d\bar{r}_p \quad \Delta \bar{r}_{o'} \rightarrow d\bar{r}_{o'} \quad (\Delta \bar{r}_p)_{xyz} \rightarrow (d\bar{r}_p)_{xyz}$$

Q: how can we write the rotation matrices?

$$R_f = R_z R_y R_x R_o$$

$R_z R_y R_x$  - sequence of body-fixed rotations from  $R_o$  to  $R_f$

If these angles are infinitesimal, the  $\theta_i \rightarrow d\theta_i$

For the rotation matrix  $\cos d\theta_i = 1$  and  $\sin d\theta_i = d\theta_i$  so

$$R_x = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & d\theta_x \\ 0 & -d\theta_x & 1 \end{bmatrix}$$

$$R_y = \begin{bmatrix} 1 & 0 & -d\theta_y \\ 0 & 1 & 0 \\ d\theta_y & 0 & 1 \end{bmatrix}$$

$$R_z = \begin{bmatrix} 1 & d\theta_z & 0 \\ -d\theta_z & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$R_z R_y R_x = \begin{bmatrix} 1 & d\theta_z & -d\theta_y \\ -d\theta_z & 1 & d\theta_x \\ d\theta_y & -d\theta_x & 1 \end{bmatrix}$$

$$d\theta_i d\theta_j \approx 0$$

We can find:

$$\begin{bmatrix} d\bar{r}_p \cdot \bar{i} \\ d\bar{r}_p \cdot \bar{j} \\ d\bar{r}_p \cdot \bar{k} \end{bmatrix} = \begin{bmatrix} d\bar{r}_{o'} \cdot \bar{i} \\ d\bar{r}_{o'} \cdot \bar{j} \\ d\bar{r}_{o'} \cdot \bar{k} \end{bmatrix} + \begin{bmatrix} (d\bar{r}_p)_{xyz} \cdot \bar{i} \\ (d\bar{r}_p)_{xyz} \cdot \bar{j} \\ (d\bar{r}_p)_{xyz} \cdot \bar{k} \end{bmatrix} + \begin{bmatrix} 0 & -d\theta_z & d\theta_y \\ d\theta_z & 0 & d\theta_x \\ -d\theta_y & d\theta_x & 0 \end{bmatrix} \begin{bmatrix} x_p \\ y_p \\ z_p \end{bmatrix}$$

no p b/c infinitesimal motion means  $(\bar{r}_p)_o = (\bar{r}_p)_f$

Matrix form of:

$$d\bar{r}_p = d\bar{r}_{o'} + (d\bar{r}_p)_{xyz} + d\bar{\theta} \times \bar{r}_{p/o'}$$

$$\text{where } d\bar{\theta} = \theta_x \bar{i} + \theta_y \bar{j} + \theta_z \bar{k}$$

infinitesimal rotation vector

## Time Derivatives (cont.)

$$d\bar{r}_P = d\bar{r}_{O'} + (d\bar{r}_P)_{xyz} + d\Theta \times \bar{r}_{P/O'}$$

We know that  $\bar{v} = \frac{d\bar{r}}{dt}$ , so dividing this equation by  $dt$  gives us  $\bar{v}_P$

$$\bar{v}_P = \bar{v}_{O'} + (\bar{v}_P)_{xyz} + \bar{\omega} \times \bar{r}_{P/O'}$$

$\bar{\omega} \equiv$  angular velocity of frame  $xyz$

$(\bar{v}_P)_{xyz} \equiv$  relative velocity - velocity in  $xyz$  frame

Note that this is equivalent to

$$\bar{v}_P = \bar{v}_{O'} + \frac{d}{dt}(\bar{r}_{P/O'}) \rightarrow \frac{d}{dt}(\bar{r}_{P/O'}) = (\bar{v}_P)_{xyz} + \bar{\omega} \times \bar{r}_{P/O'}$$

This is generally true for vectors in moving frames

$$\dot{\bar{A}} = \frac{\partial \bar{A}}{\partial t} + \bar{\omega} \times \bar{A}$$

derivative if  
the frame is  
fixed

angular velocity of  
the frame

Q: What does this mean for the derivatives of unit vectors

$$\bar{e} \rightarrow \dot{\bar{e}} = \frac{\partial \bar{e}}{\partial t} + \bar{\omega} \times \bar{e}$$

The length of  
 $\bar{e}$  is constant

## Angular Velocity and Acceleration (Sec. 3.4)

If angular velocity is  $\bar{\omega}$ , then angular acceleration is  $\bar{\alpha} = \frac{d\bar{\omega}}{dt}$

Total time deriv.

$\bar{\omega} = \sum \omega_n \bar{e}_n$  ← sum of simple rotations about axes  $\bar{e}_n$  } Define  $\bar{e}_n = \bar{i}_n, \bar{j}_n, \text{ or } \bar{k}_n$

Define  $\bar{\Omega}_n$  as the angular velocity of frame  $n$  ( $x_n y_n z_n$ )

$$\bar{\alpha} = \sum \left( \dot{\omega}_n \bar{e}_n + \omega_n \dot{\bar{e}}_n \right)$$

Q: What is  $\dot{\bar{e}}_n$ ?

$$\dot{\bar{e}}_n = \bar{\Omega}_n \times \bar{e}_n$$

So, 
$$\bar{\alpha} = \sum \left( \underbrace{\dot{\omega}_n \bar{e}_n}_{\text{change in rotation rate}} + \underbrace{\omega_n (\bar{\Omega}_n \times \bar{e}_n)}_{\text{change in orientation of } \bar{e}_n} \right)$$

change in rotation rate

change in orientation of  $\bar{e}_n$

can have angular acceleration even if  $\dot{\omega}_n = 0$  (constant rotation rate)

• This term almost never shows up in planar cases seen in undergrad dynamics

• Intuition from planar case can be wrong because it ignores this term

## Procedure to Find Angular Velocity and Acceleration (Sec 3.4.2)

- Examine rotation, write it as a series of simple rotations,  $\omega_n$
- For each  $\omega_n$ , define a frame where  $\omega_n$  aligns with  $\bar{i}_n, \bar{j}_n$ , or  $\bar{k}_n$
- Sum rotations into total angular velocity vector

$$\bar{\omega} = \omega_1 \bar{e}_1 + \omega_2 \bar{e}_2 + \dots + \omega_n \bar{e}_n$$

- Find  $\Omega_n$ , from angular velocity, using a similar method
- Solve for  $\bar{\alpha}$  by taking the full time deriv.

$$\bar{\alpha} = \underbrace{\dot{\omega}_1 \bar{e}_1}_{\text{Accel from change of } \omega_1} + \underbrace{\omega_1 (\bar{\Omega}_1 \times \bar{e}_1)}_{\text{Accel from change of } \bar{e}_1 \text{ direction}} + \underbrace{\dot{\omega}_2 \bar{e}_2}_{\text{Accel from change of } \omega_2} + \underbrace{\omega_2 (\bar{\Omega}_2 \times \bar{e}_2)}_{\text{Accel from change of } \bar{e}_2 \text{ direction}} + \dots + \underbrace{\dot{\omega}_n \bar{e}_n}_{\text{Accel from change of } \omega_n} + \underbrace{\omega_n (\bar{\Omega}_n \times \bar{e}_n)}_{\text{Accel from change of } \bar{e}_n \text{ direction}}$$

- Keys:
- must write all vectors in same frame before summing
  - choice of frames can simplify the "math" of the procedure